

STANDING SOLITARY EULER-KORTEWEG WAVES ARE UNSTABLE

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1. THE RESULT

The Euler-Korteweg system is given by the equations

$$(1) \quad \begin{aligned} V_t - U_y &= 0, \\ U_t + p(V)_y &= -(\kappa(V)V_{yy} + \frac{1}{2}(\kappa(V))_y V_y)_y, \end{aligned}$$

with $\kappa(V) > 0$. System (1), and notably its solitary waves

$$\begin{pmatrix} V \\ U \end{pmatrix}(x, t) = \begin{pmatrix} v \\ u \end{pmatrix}(x - ct) \quad \text{with} \quad \begin{pmatrix} v \\ u \end{pmatrix}(\pm\infty) = \begin{pmatrix} v_* \\ u_* \end{pmatrix},$$

have been intensely studied by Benzoni-Gavage et al. in [1, 2, 3].

Definition 1. [3] *A traveling wave (v, u) of (1) is called orbitally stable if for each $\varepsilon > 0$, there exists a $\delta > 0$ such that for any solution $(V, U) \in (v, u) + C([0, T]; H^3(\mathbb{R}) \times H^2(\mathbb{R}))$ of (1), closeness at initial time,*

$$\|(V, U)(\cdot, 0) - (v, u)(\cdot)\|_{H^1 \times L^2} < \delta$$

implies closeness at any time

$$\inf_{\sigma \in \mathbb{R}} \|(V, U)(\cdot, t) - (v, u)(\cdot + \sigma)\|_{H^1 \times L^2} < \varepsilon \quad \text{for all } t > 0.$$

The following is the point of this short note.

Theorem 1. *All non-trivial standing solitary Euler-Korteweg waves are not orbitally stable.*

2. THE PROOF

For fixed base state v_* , the solitary waves homoclinic to v_* occur in families (u^c, v^c) parametrized by their speed c . The proof of Theorem 1 is based on

Lemma 1. [1] *A solitary wave (u^{c_*}, v^{c_*}) is orbitally unstable if the moment of instability*

$$m(c) = \int_{-\infty}^{\infty} \kappa(v) v'^2 d\xi$$

is not convex at $c = c_$.*

Theorem 1 follows from

Lemma 2. $m''(0) < 0$.

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To prove Lemma 2, we recall that with

$$(2) \quad F(v, c) = -f(v) + f(v_*) - p(v_*)(v - v_*) + \frac{1}{2}c^2(v - v_*)^2, \quad -\frac{df(v)}{dv} = p(v),$$

the profile equation

$$\kappa(v)v'' + \frac{1}{2}(\kappa(v))'v' = -\frac{\partial F(v, c)}{\partial v}$$

possesses (cf. [3]) a first integral given by

$$(3) \quad I(v, v') = \frac{1}{2}\kappa(v)v'^2 + F(v, c).$$

As

$$m(c) = 2 \int_{v_*}^{v_m(c)} \kappa(v)v' dv$$

with $v_*, v_m(c) > v_*$ consecutive zeros of $F(\cdot, c)$. Since $I(v, v') \equiv 0$ along solutions, we have

$$\begin{aligned} m(c) &= 2 \int_{v_*}^{v_m(c)} (\kappa(v))^{1/2} (-2F(v, c))^{1/2} dv \\ &= 4 \int_0^{(v_m(c) - v_*)^{1/2}} (\kappa(v_m(c) - w^2))^{1/2} (-2F(v_m(c) - w^2, c))^{1/2} w dw, \end{aligned}$$

where $w := (v_m(c) - v)^{1/2}$ (cf. [6]). The first derivative of m is

$$\begin{aligned} m'(c) &= 4 \int \frac{d}{dc} \left\{ (\kappa(v_m(c) - w^2))^{1/2} (-2F(v_m(c) - w^2, c))^{1/2} \right\} w dw \\ &= 4 \int \frac{\kappa_v(v_m(c) - w^2)v'_m(c)}{2(\kappa(v_m(c) - w^2))^{1/2}} (-2F(v_m(c) - w^2, c))^{1/2} w dw + \\ &\quad + 4 \int \frac{(\kappa(v_m(c) - w^2))^{1/2} (-F_v(v_m(c) - w^2, c)v'_m(c) - F_c(v_m(c) - w^2, c))}{(-2F(v_m(c) - w^2, c))^{1/2}} w dw \end{aligned}$$

which, due to

$$\frac{\partial}{\partial v} \left((\kappa(v)(-2F(v, c)))^{1/2} \right) = \frac{\kappa_v(v)(-2F(v, c))^{1/2}}{2(\kappa(v))^{1/2}} - \frac{\kappa(v)^{1/2}F_v(v, c)}{(-2F(v, c))^{1/2}},$$

simplifies to

$$m'(c) = -4 \int \frac{(\kappa(v_m(c) - w^2))^{1/2} F_c(v_m(c) - w^2, c)}{(-2F(v_m(c) - w^2, c))^{1/2}} w dw.$$

The second derivative of m can be written in the form

$$m''(c) = 2 \int_{v_*}^{v_m(c)} \frac{A(v, c) + B(v, c)}{(\kappa(v))^{1/2} (-2F(v, c))^{3/2}} dv$$

with $A(v, c) = F(v, c)\kappa_v(v)v'_m(c)F_c(v, c)$ and

$$B(v, c) = \kappa(v) \left((v - v_*) \left(2F(v, c) \left((v - v_*) + 2cv'_m(c) \right) - c(v - v_*) \left(F_v(v, c)v'_m(c) + F_c(v, c) \right) \right) \right).$$

As $F_c(v, 0) = 0 = A(v, 0)$ and

$$B(v, 0) = 2\kappa(v)(v - v_*)^2 F(v, 0) < 0 \quad \text{for all } v \in (v_*, v_m(0)),$$

we indeed have $m''(0) < 0$.

Remark. Certain of the results by Zumbrun [7] on the Bona-Sachs model ($\kappa(V) \equiv 1$ and $p(V) = -V + V^q$ with $q \geq 2$) and by De Bouard [4] on a Gross-Pitaevskii model ($\kappa(V) = 1/(4V^4)$ and $p(V) = \alpha/V^2 - \beta/V^3$ with $\alpha, \beta > 0$) appear as interesting special cases of Theorem 1.

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REFERENCES

- [1] S. Benzoni-Gavage, *Spectral transverse instability of solitary waves in Korteweg fluids*, J. Math. Anal. Appl. **361**, 2010, 338–357.
- [2] S. Benzoni-Gavage, R. Danchin, S. Descombes, *Well-posedness of one-dimensional Korteweg Models*, Electron. J. Differential Equations **59**, 2006, 1–35.
- [3] S. Benzoni-Gavage, R. Danchin, S. Descombes, D. Jamet, *Structure of Korteweg models and stability of diffusive interfaces*, Interfaces Free Bound. **7**, 2005, 371–414.
- [4] A. De Bouard, *Instability of Stationary Bubbles*, SIAM J. Math. Anal. **26**, 1995, 566–582.
- [5] M. Grillakis, J. Shatah, W. Strauss, *Stability Theory of Solitary Waves in the Presence of Symmetry, I*, J. Funct. Anal. **74**, 1987, 160–197.
- [6] J. Höwing, *Stability of large- and small-amplitude solitary waves in the generalized Korteweg-de Vries and Euler-Korteweg/Boussinesq equations*, J. Differential Equations **251**, 2011, 2515–2533.
- [7] K. Zumbrun, *A sharp stability criterion for soliton-type propagating phase boundaries in Korteweg’s model*, Z. Anal. Anwend. **27**, 2008, 11–30.

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